

Three-Particle correlations from PHENIX to investigate the properties of the QGP at RHIC

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1 There is now a consensus of opinion that in heavy ion collisions at RHIC energies a strongly
2 interacting state of matter resembling a near-perfect liquid (termed sQGP) is formed. At-
3 tention is now being focused on elucidating the properties of this medium e.g. size, shape,
4 collectivity, viscosity, closeness to the critical point etc. In this work we describe the efforts in
5 the soft and hard sectors to quantify these properties. In particular, the study of jet-medium
6 interactions are described via correlation based studies. Details of three particle correlation
7 studies of RHIC data will be discussed. Such studies are important as they provide an avenue
8 for estimates of the viscosity and the sound speed of the medium.

Systematic Study of the Optimization Potential for Di-Lepton Measurements in the CBM Experiment

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9 CBM (Compressed Baryonic Matter) will be the dedicated heavy-ion experiment at the
10 future FAIR (Facility for Antiproton and Ion Research) complex at GSI in Darmstadt,
11 Germany. With CBM we will explore the QCD phase diagram in the region of moderate
12 temperatures but very high baryon densities with A+A collisions of 10-45 AGeV beam
13 energy. A key goal will be the search for the phase transition between hadronic and partonic
14 matter and a possible critical endpoint of this transition. Moreover, CBM will aim for an
15 understanding and the characterization of the high density baryonic matter created in these
16 A+A collisions.

17 A central observable for these studies are lepton pairs, emitted from the hot and dense phase
18 of the heavy-ion collisions. Since leptons leave the fireball without further interaction, their
19 study can provide information on in-medium properties of vector mesons or on a possible
20 restoration of the chiral symmetry.

21 Currently the design of the CBM experiment is being optimized for precision measurements
22 of these observables. In this presentation we will focus on the electron decay channel of
23 the vector mesons. The main challenge in this process is to effectively reduce the combi-
24 natorial background. Potential contributions to the background as Dalitz-decays or gamma
25 conversion have to be studied in detail and strategies have to be developed to increase the

26 signal-to-background ratio.

27 In this presentation, a systematic study on the development of an optimized detector setup
28 for a high precision di-electron measurement will be presented. This includes modifications
29 of the proposed setup of the CBM detector, the introduction of new detector components,
30 establishing of necessary purity levels and particle identification strategies.

Neutral pion production with respect to reaction plane in Au+Au collisions at RHIC-PHENIX

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31 It has been observed that the yield of neutral pions at high transverse momentum ($p_T >$
32 $5 \text{ GeV}/c$) region is strongly suppressed in central Au+Au collisions at Relativistic Heavy Ion
33 Collider (RHIC), compared to the one expected in p+p collisions. This suppression may be
34 due to an energy loss of hard scattered partons in the medium (jet quenching), that results in
35 a decrease of the yield at a given p_T . The magnitude of the suppression would depend on the
36 path length of scattering partons in the medium, and therefore is associated with azimuthal
37 angle from reaction plane in non-central collisions. Studying the path length dependence of
38 energy loss would give additional information on understanding the energy loss mechanism.
39 We discuss the parton energy loss mechanism using the nuclear modification factor (R_{AA})
40 of neutral pion with respect to reaction plane. A new reaction plane detector was installed
41 in the PHENIX detector in RHIC Year-7 run, and improved the reaction plane resolution.
42 More precise measurement of the hadron suppression with respect to path length is expected
43 using the detector. I will report about analysis status of neutral pion production in Au+Au
44 collisions at $\sqrt{s_{NN}} = 200 \text{ GeV}$.

Lattice Thermodynamics of ϕ^4 theory in 4 space-time dimensions Debasish Banerjee^a, Saumen Datta^a and Sourendu Gupta^a

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45 Perturbative calculations for QCD in the deconfined phase are problematic. A naive pertur-
46 bation expansion runs into infrared singularities, which can be taken care of by a suitable
47 resummation, taking into account the thermal mass of the electric gluon, and the fact that
48 the very long distance physics is non-perturbative. The resulting resummed perturbation
49 series, however is non-analytic in $\alpha_s = \frac{g_s^2}{4\pi}$ where g_s is the strong coupling constant and has
50 bad convergence properties at temperatures of interest to heavy-ion collision experiments.

52 Similar problems also occur in perturbative calculations of massless scalar field. A simpler
 53 system where some of the problems can be studied is a real scalar field at finite temperature.
 54 Infra-red divergences also appear in the small mass limit, and resummation schemes are
 55 called for, but, the resulting perturbation series for the thermodynamic quantities is, again,
 56 non-analytic in the coupling and possess bad convergence properties. The dependence of the
 57 thermodynamic quantities can be studied as a function of the mass parameter in the theory,
 58 especially in the limit of small mass, where the divergence structure of the theory would
 59 show up. Various resummation schemes are available in the literature for this.

60
 61 In this work, we carry out a non-perturbative evaluation of the pressure and energy density
 62 on the lattice for a real, single component ϕ^4 theory. In our work, we calculate the thermody-
 63 namic quantities for two different physical masses, each of which determine a renormalisation
 64 group trajectory in the space of the couplings. By starting from the broken phase, and staying
 65 on the RG trajectory, we determine the critical temperature, T_c , when we run into the sym-
 66 metric phase. The relevant thermodynamic quantities are evaluated at temperatures upto 3
 67 T_c . In such a limit, comparisons can be made with the results available from the perturbation
 68 theory.

SELF-CONSISTENT QUASI-PARTICLE MODEL FOR QGP

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69 Quasi-particle model (qQGP) for QGP seems to be a successful phenomenological model
 70 to explain the non-ideal behaviour seen in the lattice simulation of QCD (LGT) and the
 71 relativistic heavy ion collisions (RHICs) experiments. Here, one replaces the system with
 72 real quarks and gluons, interacting via QCD, by non-interacting quasi-partons, having ad-
 73 ditional thermal masses. There are verities of qQGP models [1] with few system dependent
 74 adjustable parameters. Recently, we proposed a new formulation of qQGP, with a **single**
 75 system dependent adjustable parameter, which explains very well the LGT results on pure
 76 gauge, 2f, 3f and (2+1)f QGP systems. In many works, the thermal masses were taken from
 77 the finite temperature QCD perturbative calculations which is based on the temperature
 78 dependent propagators appropriate to ideal system and valid for very high temperature.
 79 But the system is **non-ideal** and also **not** very high temperature so that the perturbation
 80 theory is valid. Here we present and compare 3 models based on our new formulation [2] of
 81 qQGP which does not need the reformulation of statistical mechanics (SM) [3], but use the
 82 standard SM where one generally obtain the energy density and number density from the
 83 definition of grand canonical ensemble and derive all other thermodynamic (TD) functions,
 84 including the pressure without any TD inconsistency. Model-I is based on the perturbative
 85 thermal masses usually used in all qQGP calculations. In Model-II, we assume that ther-
 86 mal masses depend on density (not ideal gas) which itself is a functional of thermal masses
 87 and one need to solve the problem self-consistently. But we assume the same approximate
 88 dispersion relation (ω Vs k) used in other qQGP which is appropriate for $T \gg T_c$, which
 89 is again questionable. Hence, in Model-III, we redo our self-consistent qQGP, but with the

90 exact dispersion relation obtained in QCD perturbative and resummation methods. We see
91 that, for all systems (pure gauge, 2f, 3f and (2+1)f), results improve as we go from Model-I
92 to Model-III. At high temperature, all 3 models matches but differ as we go near $T = T_c$.
93 All 3 models have only one system dependent adjustable parameter and explain surprisingly
94 well the LGT results.

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Measurement of High p_T Identified Charged Hadron Spectra in Au+Au $\sqrt{s_{NN}} = 200$ GeV Collisions by the PHENIX Experiment

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102 In Run-7 of RHIC operations the PHENIX experiment collected over 5 billion minimum bias
103 Au+Au $\sqrt{s_{NN}} = 200$ GeV events. The large dataset and new detector capabilities greatly
104 extend the transverse momentum (p_T) range of identified charged hadron measurements in
105 PHENIX. The Time-of-Flight West detector, one of the new detectors used in PHENIX
106 during the latest run, allows for seamless track-by-track particle identification of $\pi/K/p/d$
107 up to $p_T \approx 9$ GeV/c when used in conjunction with the Aerogel Cherenkov Counters. The
108 nuclear modification factors (central-to-peripheral ratios) for $\pi/K/p/d$ and baryon/meson
109 ratios will be presented. These measurements, coupled with two-particle correlations and
110 flow measurements, can contribute to a more detailed understanding of the partonic energy
111 loss mechanism and the baryon anomaly observed at RHIC.

SEARCHING FOR THE COLOR GLASS CONDENSATE WITH THE STAR FMS

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112 The first runs at the RHIC collider operated at $\sqrt{s_{NN}} = 200$ GeV have demonstrated the

113 possibility to probe the gluon distribution in heavy nuclei using forward electromagnetic
114 calorimeters. In particular, early results from the Solenoidal Tracker At RHIC (STAR) using
115 prototype calorimeters (FPD/FPD++) have shown general agreement of inclusive yields of
116 forward π^0 mesons from p+p collisions at $\sqrt{s} = 200\text{GeV}$ with NLO pQCD calculations. In
117 d+Au collisions, on the other hand, a stronger suppression of forward π^0 yield is observed
118 than would be expected from shadowing effects. Exploratory measurements of azimuthal
119 correlations of the forward π^0 with charged hadrons at $\eta \sim 0$ show a recoil peak in p+p that
120 is suppressed in d+Au at low pion energy. It is shown here how these observations can be
121 qualitatively interpreted using the Color Glass Condensate (CGC) model, an effective field-
122 theory for the saturation behavior of the small-x component of the nuclear wave function.
123 A new detector, the Forward Meson Spectrometer, has been built to measure forward pion
124 and photon production to quantify gluon suppression at low x in heavy nuclei over a large
125 range in x and pt. Using CGC calculations, we show that the detector probes x down to
126 $x \approx 10^{-4}$ for inclusive particle production at $\eta \sim 4$ at $\sqrt{s_{NN}} = 200\text{GeV}$. This is well into
127 the x-range where gluon saturation effects are expected to set in. The wide electromagnetic
128 coverage ($2.5 < \eta < 4.0$) will play a crucial role in mapping the boundaries for signatures of
129 saturation (CGC) effects. The granularity of the FMS will allow discrimination of prompt
130 photons from photons produced by neutral meson decays at large rapidity.

Matter induced charge symmetry breaking

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131 The isospin symmetry is only an approximate symmetry of strong interaction. In nature,
132 for example, light quark masses are non-degenerate. At the hadronic level a mass splitting
133 is observed between the proton (p) and neutron (n) or between the various charge states
134 of the isospin meson multiplets. At the QCD level the symmetry breaking effects can be
135 parametrized in terms of a small parameter $\epsilon = (m_d - m_u)/(m_u + m_d)$ [1]. Consequently,
136 in the low energy limit of QCD, mixing of various isospin resonance states like ρ - ω mixing
137 or π - η mixing can be expressed in terms of similar parameter involving n and p masses
138 [2]. This, clearly, is an explicit symmetry breaking by the Hamiltonian itself. Such mixings
139 of various isospin states are responsible for charge symmetry violation (CSV) in nucleon-
140 nucleon (NN) interactions. In fact, there exist lot of calculations directed towards the
141 construction of various CSV NN interactions which can be classified as class III and IV
142 potentials. In addition, even the mass difference of various charged states of mesons, like π^0 ,
143 π^\pm can break isospin symmetry which is responsible for the breaking of charge independence
144 of NN interaction even if nucleon masses are taken to be degenerate.

145 In the present work, we consider another mechanism of such symmetry breaking phenomenon
146 which unlike the explicit symmetry breaking term in the Hamiltonian, involves the ground
147 state of nuclear matter containing unequal densities of neutron and proton. This is akin to
148 the spontaneous symmetry breaking (SSB) where the Hamiltonian respects the symmetry
149 but it is broken implicitly by the ground state. Clearly this is a density induced effect. The
150 possibility of such mixing was discussed first in ref.[3] and subsequently in [4]. However all
151 these calculations were limited in the time-like domain, while we extend to formalisms to
152 the space-like region and construct density dependent CSV potential.

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Energy Loss of Heavy Quarks from Single Muon Spectra in Pb-Pb collisions at $\sqrt{s_{NN}} = 5.5$ TeV

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158 We will present generator-level studies of the effect of radiative energy loss on heavy quarks
 159 (charm and beauty) produced in central heavy-ion (Pb-Pb) collisions at $\sqrt{s_{NN}} = 5.5$ TeV.
 160 We use PYTHIA and HYDJET Monte Carlos to obtain the nuclear modification factor
 161 $R_{AA}(p_T)$ of single muons from semi-leptonic decays of heavy-flavoured mesons (D and B),
 162 in the $|\eta^\mu| < 2.5$, $p_{tot}^\mu > 5$ GeV/c kinematic range (typical of LHC detectors like ATLAS or
 163 CMS), and discuss its relevance as a signature for medium-induced radiated parton energy
 164 loss on charm and bottom quarks. Based on the different suppression levels exhibited by
 165 beauty, charm and light-quark channels at intermediate and high p_T , a semi-quantitative
 166 dependence of the amount of energy loss on quark mass and flavour can be established from
 167 the single inclusive muon spectra.

Measurement of the D^0 meson Production through Hadronic Decay Channel at PHENIX

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168 The Open Charm production have always been considered as one of the unique and unbiased
 169 probes to study the properties of the dense matter produced in Heavy Ion collisions at RHIC.
 170 PHENIX experiment at RHIC has done a tremendous job of measuring the production rate of
 171 the Heavy Flavor particles through their semi-leptonic decay channels. By studying the pro-
 172 duction in a different colliding systems (p+p, d+Au, Au+Au) we were able to separate both
 173 initial and final state effects in nucleon+nucleon collisions at 200 GeV/N. The unanswered
 174 question for those measurements was a relative contribution of Open Charm to the mixture
 175 of Open Charm and Open Bottom particle decays that were accessible experimentally in
 176 semi-leptonic channel. The results of the direct measurement of D^0 meson production in
 177 proton+proton collisions through the hadronic decay channel will be presented. This mea-

178 surement will should determine the contribution of the Open Charm to the mix and also
179 constrain the theoretical models for the Open Charm production at RHIC energies. The
180 current status of analysis, as well as comparison with the most resent theoretical predictions
181 will be presented.

Decomposition of Awayside Components of Dijet Correlations in Au+Au Collisions at $\sqrt{s_{NN}} = 200$ GeV at PHENIX

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182 Di-hadron angular correlations offer a powerful method to study the interactions of jets with
183 the hot, dense medium in heavy ion collisions. Since fragments from opposing jets typically
184 form back-to-back pairs, the dihadron angular correlation can also be used to probe the
185 medium response of the energy deposited by parton energy loss. At intermediate transverse
186 momentum (p_T), dijet correlations have a local minimum at $\Delta\phi = \pi$ in Au+Au. The
187 modified shape suggests a medium response to energy deposited by the transiting parton.

188 We present inclusive photon-hadron angular correlations from $\sqrt{s_{NN}} = 200$ GeV Au+Au
189 collisions in PHENIX. At this intermediate p_T (below 5 GeV/c), these correlations are ap-
190 proximately all from meson-hadron pairs. We decompose the awayside region in the two
191 particle correlation into "head", corresponding to jets which punch through the other side
192 of the medium, and "shoulder", which contains the medium response. We present the per
193 trigger yield in each component and study the transverse momentum and energy balance
194 with the nearside which may also be modified. Au+Au is compared with $\sqrt{s_{NN}} = 200$ GeV
195 p+p collision measured by PHENIX.

PHENIX measurement of centrality dependent charged hadron production in deuteron-gold and nucleon-gold collisions at $\sqrt{s_{NN}} = 200$ GeV

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196 We present transverse momentum (p_T) spectra of charged hadrons measured in deuteron-gold
197 and nucleon-gold collisions in the PHENIX experiment at $\sqrt{s_{NN}} = 200$ for four centrality
198 classes. Nucleon-gold collisions were selected by tagging events in which a spectator nu-
199 cleon was observed in one of two forward rapidity detectors. The spectra and yields were
200 investigated as a function of the number of binary nucleon-nucleon collisions, ν , suffered by
201 deuteron nucleons. A comparison of charged particle yields to those in $p + p$ collisions show
202 that yield per nucleon-nucleon collision saturates with ν for high momentum particles. We

203 also present the charged hadron to neutral pion ratios as a function of p_T .

Determination of the effect of final state energy loss on the acoplanarity and momentum imbalance in jets

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204 The energy loss of partons leaving the hot media created in the heavy ion collisions has been
205 at the center of the interest of the experimental and theoretical communities of the field.
206 Hence an increase in the transverse momentum imbalance and acoplanarity of dijets may
207 be a way to study the parton energy loss [1]. In order to study the effect, we have made
208 simulations of dijet events where the intrinsic k_t of the partons, the initial and final state
209 radiation are taken into account. The results of these simulations as a function of the center
210 of mass energy of the collision will be presented.

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NEUTRON-ANTINEUTRON OSCILLATION AND ITS IMPLICATION ON A NEUTRON STAR

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The neutron-antineutron ($n\bar{n}$) oscillation is the phenomenon when neutrons change spontaneously to antineutrons and vice versa with an oscillation time $\tau_{n\bar{n}}$. If this process takes place in nature then we have Baryon number oscillation (violation) process with $\Delta B = 2$.

$$n(B = +1) \leftrightarrow \bar{n}(B = -1)$$

In vacuum, the probability of $n\bar{n}$ transition can be calculated to be

$$P(n(t) = \bar{n}) = |\langle \bar{n} | n(t) \rangle|^2 = \sin^2 \delta t = \sin^2 \frac{t}{\tau_{n\bar{n}}}$$

where t denotes the time. In the presence of magnetic field, this probability can be deduced as

$$P(n(t) = \bar{n}) = \sin^2 2\theta \sin^2 \left(\frac{\Delta E \tau_{n\bar{n}}}{2} \right)$$

where

$$\frac{\Delta E}{2} = \sqrt{(\vec{\mu}_n \cdot \vec{B})^2 + \delta^2}$$

and

$$\tan \theta = \frac{\delta}{\left(\vec{\mu}_n \cdot \vec{B} + \sqrt{(\vec{\mu}_n \cdot \vec{B})^2 + \delta^2} \right)} .$$

214 In the above B is the magnetic field strength, $\vec{\mu}_n$ being the neutron magnetic moment and the
215 transition amplitude $\delta = \frac{1}{\tau_{n\bar{n}}}$. The oscillation probability in matter can also be calculated.
216 Neutron stars are dense collapsed core of some massive stars with typical mass of $1.4M_{\odot}$ and
217 radius ~ 1.4 Km with a magnetic field around 10^{12} Gauss. At the initial stage the power
218 loss of the neutron star is through the URCA process like $n + n \rightarrow n + p + e + \nu$ producing
219 neutrinos. At a density around 4×10^{11} gm/cc inside the neutron star the neutrons are
220 abundant. The $n\bar{n}$ oscillation, if significant may affect these processes and the stability of
221 a neutron star. In this work, the effect, if any, of such oscillation on neutron star will be
222 investigated.

ON THE NEW POSSIBILITY TO DESCRIBE THE TIME EVOLUTION OF THE CHEMICAL POTENTIAL IN NUCLEAR COLLISIONS AT HIGH ENERGIES

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223 At the last Quark Matter Meeting (Shanghai 2006 - China) we proposed the introduction of
224 an empirical scale of time for nuclear reactions at high energies, inspired by the HUBBLE
225 Cosmology. The goal of this work is to determine the time evolution of the chemical potentials
226 for charged particles (p^{\pm} , K^{\pm} , π^{\pm}) like those obtained in Au+Au at 200 GeV using the
227 blast-wave model.

Studying the sources of high p_T direct photons in Au+Au collisions with the PHENIX detector at RHIC

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228 In addition to primordial hard scattering at least three other mechanisms are assumed to
229 contribute to the observed high p_T (>5 GeV/c) direct photon spectrum in relativistic heavy

230 ion collisions: fragmentation, jet-photon conversion, and Bremsstrahlung. The azimuthal
231 asymmetries are predicted to be different for these mechanisms, notably jet-conversion and
232 Bremsstrahlung photons should exhibit negative “flow”. Also, hard scattering and jet-
233 conversion produces photons without any accompanying jet particles, although such “iso-
234 lation” is blurred by the underlying Au+Au event. We will review current results of high
235 p_T direct photon production at different energies, then using the new, high resolution reac-
236 tion plane detector of PHENIX and loose isolation criteria we will study the possibility of
237 statistically separating photons from the four basic production sources listed above.

238 **High- p_T Measurement of Azimuthal Anisotropy of Electrons from**
239 **Semi-leptonic Decay of Open Heavy Flavor Mesons in Au+Au Collisions at**
240 **$\sqrt{s_{NN}} = 200\text{GeV}$ by the PHENIX Experiment**

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241 The azimuthal anisotropy of open heavy flavor mesons at RHIC provides information about
242 the early stages of heavy ion collisions. PHENIX observed a large azimuthal anisotropy
243 parameter v_2 of electrons from heavy flavor decays from RHIC Run4 at low transverse mo-
244 mentum. Together with the observation of constituent quark scaling of v_2 for light hadrons,
245 the electron v_2 is consistent in coalescence models with charm quarks flowing as much as
246 light quarks. With the addition of a reaction plane detector in PHENIX and the increased
247 statistics of RHIC Run7, v_2 of electrons can be measured with much-improved precision,
248 answering the question of whether heavy flavor continues to flow in the range of p_T in which
249 beauty production is believed to be comparable to that of charm. We studied the p_T and
250 collision centrality dependency of the azimuthal anisotropy of single electrons from open
251 heavy flavor decays from RHIC Run7 data in PHENIX.

252 **Measurement of excited charmonium state and study of J/ψ polarization in**
253 **PHENIX experiment at RHIC**

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254 Determining the J/ψ polarization will provide constraint on charmonium formation mech-
255 anisms. Measuring the feed down of ψ' to J/ψ production is one of the main ingredients

256 to obtain information on J/ψ prompt production as well as to understand its anomalous
257 suppression observed in AA collisions at RHIC.

258 The PHENIX Collaboration has studied charmonium production in $p+p$ collisions at RHIC
259 using the sample collected during the 2006 data taking period through its decays into the
260 dielectron channel at mid rapidity. This high statistics data set allowed us to study polar-
261 ization of the J/ψ and the ψ' to J/ψ ratio at RHIC energies.

A first study of the dimuon mass continuum components in p+p collisions at $\sqrt{s} = 200$ GeV

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262 The invariant mass distribution of the dimuons produced in heavy ion collisions contains
263 resonances, especially J/ψ that has been extensively used as a probe of the Quark Gluon
264 Plasma (QGP). In addition, the dimuon mass distribution above $2 \text{ GeV}/c^2$ includes con-
265 tinuum components such as Drell-Yan, open charm and open beauty. Heavy flavours carry
266 information on the medium properties in which they have propagated and thus a study of the
267 continuum in p+p, d+A and A+A collisions provides complementary signals to understand
268 detailed aspects of matter produced in these collisions. Extracting the continuum compo-
269 nents yields from a complex spectrum cannot be done without any knowledge of their mass
270 distributions (line shapes). In a first step, a full PHENIX simulation of the three continuum
271 components is done. Events produced by Pythia are sent into the PHENIX simulation and
272 reconstruction software. The vertex detector resolution was tuned to reproduce the J/ψ
273 width measured by PHENIX in p+p collisions at 200 GeV. The relative magnitudes of the
274 components are then obtained by fitting the experimental p+p data. Result from this study
275 will be presented.

Measurement of Electron-Muon Correlations from Semi-Leptonic D Decay in 200 GeV p+p Collisions at RHIC-PHENIX

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276 Charm production is a valuable probe of the early stages of a heavy ion collision. Corre-
277 lated electron-muon pairs are a signature of semi-leptonic D decays, and a measurement of
278 D mesons provides information on charm quark energy loss in the hot medium. The energy
279 loss of heavy quarks is still not fully understood, so it is vital to investigate different decay
280 channels of charm mesons to better understand this process. Measurements of electron-muon
281 pairs suffer less from background than e^+e^- or mu^+mu^- pairs since neither direct lepton
282 production nor resonance decays produce this type of correlated signal. Another advantage

283 results from measuring the electrons in the PHENIX central arms and muons in the forward
284 direction, since this probes open charm in a rapidity region different from previous measure-
285 ments. Studying electron-muon pairs in p+p collisions provides an important baseline for
286 the study of these processes in d+Au and Au+Au collisions. The data in this analysis was
287 obtained during the 2006 RHIC run of p+p collisions at 200 GeV. The current status of this
288 analysis will be presented.

Probing RHIC Matter with direct-photon Jet and Jet-Jet Correlations in PHENIX^a

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289 2-particle correlations from Direct Photon - Jet and Jet-Jet Pairs are important for studying
290 Energy Loss of Hard Probes in the matter created at RHIC. We present details of the Direct
291 Photon-h and pi0-h analyses from p+p and A+A datasets ranging from Run4 to the newest
292 Run7.

Analysis of features of multiplicity distributions of charged hadrons at RHIC

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293 It is found that Collective Flow Model(CFM)which can successfully analyze charged particle
294 distributions at AGS and lower SPS (less than 20Gev/n) , fails to analyze that of RHIC. The
295 tail of distribution of charged particle at RHIC has a jump away from the collective model
296 calculation as the energy increase. Thermalization Component Model (TCM) is presented
297 basing on collective flow to study the multiplicity distributions at RHIC in this paper. It
298 is realized that the limitation of phase space of collective flow can denote that of thermalization
299 region. By comparing the contributions of particle productions from thermalization region at
300 different energies and different centrality, we can deep our study on the feature of collective
301 movement at RHIC.

302 Detailed analysis of thermalization relation with centrality and energies at RHIC is presented
303 in this paper. We develop CFM to TCM to discuss the thermalization limitations and
304 contributions at different energy regions and different centrality in this paper.

305 Collective flow in heavy-ion collisions is an unavoidable consequence of thermalization. It is
306 realized that the limitation of phase space of collective flow can denote that of thermalization

307 region. We have shown that the Thermalization Component Model gives results consistent
308 with the experimental data for particle production at AGS, SPS and RHIC. Especially we
309 present a detail discussion of Au-Au and Cu-Cu collisions at different centrality at $\sqrt{s} = 62.4$
310 and $\sqrt{s} = 200$ GeV. The percentage ratios of contributions of particle production from the
311 thermalization region are the largest at AGS, and decrease as collision energies decrease at
312 SPS and RHIC, but seem to reach saturation when $\sqrt{s}=62.4 - 200$ GeV. It is also found that
313 the limitation of the flow show a relation of straight line with $\ln\sqrt{(s)}$. We can predict the
314 thermalization limitation at future LHC experimental data.

315 It is shown from our study that the percentage ratios of particle production from thermaliza-
316 tion regions increase as the increase of the centrality at RHIC. We also study the percentage
317 ratios of contribution from thermalization regions by compare two collision systems(Au+Au
318 with Cu+Cu) and two collision energies at RHIC. It is found that the contribution ratio is
319 larger for the smaller collision system (Cu+Cu) at $\sqrt{s} = 62.4$ GeV , but shows independent
320 of the collision system at $\sqrt{s} = 200$ GeV .

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STAR Multi-Strange Hyperon Elliptic Flow and Parton Collectivity from 200GeV Au+Au and Cu+Cu Collisions

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324 We will present measurements of elliptic flow (v_2) for multi-strange hyperons (Omega and
325 Cascade) from 200GeV Au+Au and Cu+Cu collisions. Both the signal to background ratio
326 and v_2 are calculated as a function of invariant mass. For a given pt bin, the observed v_2
327 can be treated as a superposition of background v_2 and omega v_2 . A fit can be made to
328 the v_2 vs. mass distribution and the hyperon v_2 extracted. This method has successfully
329 been used for STAR phi meson v_2 measurements[1]. The physics implications on coalescence
330 models and parton collectivity will be discussed.

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Heavy Quarkonia Production in Heavy Ion Collisions at PHENIX

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333 It has been suggested that heavy quarkonia suppression could be a strong signature of the
334 existence of QGP color screening effects in high-energy heavy ion collisions. At RHIC, we
335 have observed significant suppressions of J/Ψ production in both CuCu and AuAu collisions
336 relative to the yield in pp system. However, the current model calculations of J/Ψ suppression
337 include several competing physics processes, and it is very challenging to distinguish various
338 suppression scenarios from J/Ψ measurement alone. It is expected that the weakly bound
339 Ψ' should melt at a lower temperature and thus should be suppressed much more compared
340 with J/Ψ . On the other hand, the Υ states which are much more tightly bounded compared
341 with J/Ψ , should melt at higher temperature and hence expected to be less suppressed. So
342 the measurements of higher mass heavy quarkonia states, such as Ψ' and Υ , together with
343 J/Ψ , will help us to further constrain various quarkonium suppression models in heavy ion
344 collisions. In this analysis, we study the production rates of Ψ' s and Υ s relative to J/Ψ yields
345 in pp and CuCu systems through the dimuon channel. The current status of this study will
346 be presented.

347 **Direct measurement of fragmentation photons in p+p collisions at $\sqrt{s_{NN}} = 200$**
348 **GeV with the PHENIX experiment**

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349 In heavy ion collisions direct photons have proven to be an important observable because
350 they are penetrating, and therefore remain largely insensitive to the final state effects that
351 lead to jet quenching. Perturbative QCD calculations predict a contribution to the direct
352 photon yield of up to 30% from photons produced through parton fragmentation. In heavy
353 ion collisions it is expected that this contribution is modified due to additional stimulated
354 photon bremsstrahlung as well as energy loss of the partons through gluon radiation prior to
355 fragmentation. A measure of such bremsstrahlung photons would provide direct observation
356 of the energy loss of jets in the medium. Thus measurements of the fragmentation component
357 to direct photon yields in both p+p and Au+Au collisions provides both an important test
358 of pQCD predictions and of predictions for the nuclear modification factor. A natural way of
359 picking out photons produced through jet fragmentation is to select photons associated with
360 jets on the near side using hadron-photon correlations. However, this signal is small compared
361 to the various contributions to the inclusive conditional yield from photons produced through
362 π^0 , η , and other mesonic decays, thereby requiring a very precise determination of these
363 backgrounds. We present new results from the application of this approach to PHENIX p+p
364 data at $\sqrt{s_{NN}} = 200$ GeV and discuss its potential for signal extraction in heavy ion
365 collisions.

Event-by-event fluctuations of identified particles as measured by PHENIX

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366 Phase transitions are normally associated with fluctuations on many length scales. Thus,
367 fluctuations have been proposed as a signature of the transition of a quark-gluon plasma
368 to normal hadronic matter. These fluctuations may appear as correlations between the ob-
369 served particles. Studies of these correlations could reveal properties of the phase transition
370 and the hadronization process. We present a systematic study of kaon-to-pion fluctuations
371 for Au+Au collisions at $\sqrt{s_{NN}} = 62.4$ and 200 GeV, as a function of centrality and of the mo-
372 mentum of identified pions and kaons as measured by PHENIX. Proton-to-pion fluctuations
373 are also studied. The robustness of the analysis method will be discussed and comparisons
374 will be made with different fluctuation models using Monte Carlo techniques.

Searching for strong parity violation signals in non-central Au+Au collisions at $\sqrt{s_{NN}} = 200$ GeV in PHENIX experiment at RHIC

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375 The participants in non-central Au+Au collisions could have a large orbital angular mo-
376 mentum perpendicular to the reaction plane as determined by spectators. According to a
377 recent prediction [1], it is suggested that a strong parity violation signal could be revealed
378 via preferential emission of charged particles in the direction along the orbital angular mo-
379 mentum. This would mean that positive and negative pions will be produced asymmetrically
380 with respect to the angular momentum. There would be more positive than negative pions
381 in the direction of angular momentum and more negative than positive pions in the oppo-
382 site direction of angular momentum. We are looking for such a parity violation effect with
383 charged hadron production in the PHENIX muon spectrometer with respect to the reaction
384 plane determined from the directed flow measurement. We present the current status of this
385 analysis.

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Open charm elliptic flow from single muon in Au+Au collisions at $\sqrt{s_{NN}} = 200$ GeV in PHENIX experiment at RHIC

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388 Event anisotropy is expected to have sensitivity to the early stage of ultra-relativistic heavy-
389 ion collisions at RHIC. The possible formation of a quark gluon plasma (QGP) could affect
390 how the initial anisotropy in the space coordinate is transferred into the momentum space
391 for the final state. We perform a quantitative analysis of the azimuthal distribution of single
392 muons with respect to the reaction plane with PHENIX muon spectrometer in Au+Au
393 collisions at $\sqrt{s_{NN}} = 200$ GeV. We present the current status of the open charm elliptic
394 flow measurement after subtracting decay muon and punch-through contributions from the
395 inclusive single muon flow at forward rapidity $1.5 < |\eta| < 1.8$.

396 **”Measurement of Azimuthal Anisotropy with the New Reaction Plane Detector**
397 **in the PHENIX experiment”**

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398 Azimuthal anisotropy of particle emission with respect to the reaction plane is one of the
399 most important global observables in relativistic nucleus-nucleus collisions at RHIC. The
400 observation of a large anisotropy which follows specific scaling relations over a broad range
401 of particle species is considered as evidence for the probable formation of a hot and dense
402 partonic matter in Au+Au collisions at RHIC. The accurate measurement of the reaction
403 plane is a key factor in the study of azimuthal anisotropy. In order to improve the resolution
404 of such measurements in the PHENIX experiment at RHIC, we designed and fabricated a
405 new Reaction Plane Detector (RxNP). RxNP worked very well during the PHENIX Run7
406 period and demonstrated the design performance. As a result of the upgrade, the reaction
407 plane resolution was improved by a factor of two. This allows us to improve the precision of
408 the measurements of azimuthal anisotropy for particles with high transverse momentum $p_T >$
409 4 GeV/c and to perform a more detailed study of rare particles, such as direct photons, non-
410 photonic electrons from D-meson decay, and J/psi. First experimental results of the study
411 of azimuthal anisotropy with the new PHENIX Reaction Plane Detector will be presented
412 and discussed.

INSTABILITIES OF THE GLASMA

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413 Understanding the mechanism of thermalization of hot QCD matter is one of the most
414 important problems in the physics of heavy ion collisions [1]. Since the standard scenario
415 based on perturbative hard scattering has been challenged by very short thermalization time
416 $t < 0.6$ fm suggested by the successful ideal hydrodynamical simulation for the RHIC data.
417 Some other scenarios are now extensively investigated. One of the possible scenarios will be
418 the plasma (Weibel) instability where the interaction between hard particles and soft gauge
419 field induces rapid enhancement of the soft field. The plasma instability certainly helps
420 to make the matter isotropic, but still requires long time to reach thermalization. In this
421 contribution, we discuss another scenario for the early thermalization.
422 The initial condition of the heavy ion collision is given by the Color Glass Condensate
423 (CGC), which is a weakly-coupled strong gauge field in each colliding nucleus. The matter
424 just after the collision is still well-described by such coherent Yang-Mills field. This is called
425 the glasma. If there is some instabilities in the glasma, it will help to make the matter
426 isotropic/thermalized. In fact, it is already found in the lattice numerical simulation that
427 there is instability in the glasma [2]. We analytically investigate the classical Yang-Mills
428 equation for an expanding glasma in the rapidity (η) - proper time (τ) coordinates, and
429 find there is indeed instabilities for a very simple configuration of the gauge fields (constant
430 electric field or magnetic field in the beam (z) direction). This is essentially equivalent to the
431 known instabilities in the constant Yang-Mills field in the ordinary space-time coordinates
432 [3]. Therefore, we find that the instability of electric fields is related to the Schwinger
433 mechanism, and that of magnetic fields corresponds to the Nielsen-Olesen instability. We
434 further discuss more realistic case with the CGC initial condition. The CGC is characterized
435 by the gauge field which is randomly distributed on the transverse space with the coherent
436 size being $1/Q_s$. We discuss whether the instability occurs when the transverse extent of the
437 constant gauge field is finite and of the order of $1/Q_s$. We also discuss possible relationship
438 between the glasma instability and the Weibel instability. To clarify it, we need to examine
439 the effects of back reaction in the glasma instability.

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PHENIX measurement of dihadron correlation in Au+Au collision at 200 GeV: jet quenching and medium response

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446 We present a detailed survey of the trigger p_T , partner p_T and centrality dependence of
447 near- and away-side jet shapes and yields in Au+Au collisions at $\sqrt{s_{NN}} = 200$ GeV from
448 RUN4. These measurements provide a broad overview of the different physical features that
449 come into play for different p_T ranges. In particular, the results can be interpreted as the

450 interplay between the jet fragmentation and response of the medium to quenched jets. The
451 former dominates the high p_T region, while the latter dominates the low and intermediate p_T
452 region. These results allow a detailed comparison of the similarities as well as the differences
453 of the correlation pattern between the near- and away-side, and provide new insights into
454 the physical processes of the jet-medium interactions. Together with inclusive hadron pro-
455 duction, they also allow us to quantify the role of jets at intermediate p_T , where the particle
456 production was believed to be dominated by the soft processes such as hydrodynamical flow
457 and recombination.

PHENIX measurement of dihadron $\Delta\eta$ correlation in Au+Au collision at 200 GeV

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458 The observation of a long range $\Delta\eta$ correlation at the near-side has led to many speculations
459 of the underlying physical processes responsible for such novel feature. This include several
460 models advocating the heating of medium by the jets emitted close to the surface. PHENIX
461 can access this so called "ridge" phenomena in a limited pseudo-rapidity of $|\Delta\eta| < 0.7$ in
462 Au+Au collisions. Status of this study and the comparison to $\Delta\phi$ correlation results is
463 presented.

A new type of Time-Of-Propagation (TOP) Cherenkov detector for particle identification

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464 The Time-of-Propagation (TOP) Cherenkov detector is a novel technical extension of
465 the Detection of Internally Reflected Cherenkov (DIRC) light for particle identification. In
466 contrast to the BarBar DIRC [1,2], which provides the 2-D space measurement, it measures
467 1-D time of propagation and 1-D space information. Thus TOP has the advantages of
468 being simple in construction and compact in size as well as its quick time response. It is
469 proposed to be used in the internal target experiment on the Cooling Storage Ring (CSR) at
470 Lanzhou, China [3]. In this presentation, Geant4 toolkit is used to simulate the propagation
471 of Cherenkov photon in thin quartz bar radiator. We estimate various contributions to the
472 time uncertainty, such as the wavelength dependence of the refractive index, the hit position
473 resolution and TTS of the phototube, timing jitter due to the start signal determination,
474 the quartz thickness and the charged particle tracking accuracy [4].

475 In order to estimate the PID performance, the separability (S_0) is defined as π and K for
476 example. Its dependence on the particle momentum, incident angle and propagation length,
477 are discussed respectively. From Geant4 simulation, we obtain the probability function of
478 the deviation of the time of propagation from correct/incorrect expectation for π and K
479 tracks. A log-likelihood ($\ln L$) is thus defined and calculated for incoming π and K tracks

480 by using this probability function. The clear separation of π and K tracks illustrates the
481 high PID ability of this novel TOP technique.

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r-MODES AND VISCOSITY OF SUPERFLUID QUARK STARS

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487 We present results from a study of the r-mode instability in compact stars containing color
488 superconducting quark matter. The r-mode is a toroidal pulsational instability that arises
489 for all such rapidly rotating compact stars, causing energy loss from the star in the form
490 of gravitational waves that may be detected by interferometric experiments on Earth. The
491 r-mode instability is damped by the viscosity of quark matter, thereby suppressing the
492 amplitude of the gravitational waves while allowing such compact stars to spin faster than
493 expected.

494 For stars containing quark matter in the color-flavor-locked phase, we present results for
495 the r-mode frequency, the damping timescale associated to the bulk viscosity, as well as
496 the critical rotation frequency of these stars. Employing recent results on the bulk and
497 shear viscosities of superfluid quark matter, we show that the r-mode instability is active
498 in a window of temperature that is typical of observed compact stars, suggesting that they
499 may be promising candidates for detection of color superconducting quark matter through
500 gravitational waves.

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Measurement of $\omega/\phi \rightarrow e^+e^-$ in proton + proton collisions at $\sqrt{s} = 200\text{GeV}$ at RHIC-PHENIX

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503 The light hadron masses are generated due to the spontaneous breaking of the chiral symme-
504 try. Due to the effect of chiral symmetry restoration, the mass of the light vector meson(ρ ,
505 ω , ϕ) may be shifted and/or modified in the hot matter created by high energy heavy ion
506 collisions at RHIC. The study of mass modification is an important topic to understand the
507 mechanism of generation of hadron mass.

508 Unlike hadrons, electrons do not interact strongly with the medium. The measurement
509 of electron pairs from vector meson is therefore a good probe to study chiral symmetry
510 restoration since electrons carry the original information.

511 The study of the light vector mesons in proton+proton collisions is an important baseline
512 for the various heavy ions collisions such as Au+Au, Cu+Cu and d+Au. The PHENIX
513 experiment measures ω/ϕ meson via their di-electron decays in p+p collisions at $\sqrt{s} =$
514 200GeV. We will present and discuss the result.

515 **Two- and three particle correlations of high- p_t charged** 516 **hadrons in heavy ion collisions at top SPS engery**

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517 The analysis of azimuthal correlations of hadrons with high transverse momenta provide a
518 unique tool to study the interactions of partons with the medium in an early stage of heavy
519 ion collisions.

520 A detailed analysis of two particle correlations in Pb-Au collisions from the CERES experi-
521 ment at top SPS energy will be presented for different centralities of the collisions.

522 The key observation is a non-Gaussian shape on the "away-side" of the two particle correla-
523 tion function in central collisions indicating significant interactions of partons traversing the
524 medium even at SPS energy. A study of the correlations for different charge combinations of
525 trigger and associate particles reveals charge ordering in the fragmentation process and sen-
526 sitivity to the charges of the interacting partons. Hence the analysis may reveal properties
527 of the initial as well as the final state of the collision.

528 To further investigate the pattern of the two particle correlation function an analysis based
529 on three particle correlations is presented. Different scenarios like elastic scattering of the
530 initial partons or the evolution of a mach cone can lead to the same observable shape of
531 the two particle correlation function. The three particle analysis gives strong indications
532 for a cone like emission of hadrons emerging from hard parton-parton interactions in the
533 collisions.

Beam Energy and System Size Dependence of Radial Flow Measured at PHENIX
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534 Hadron production in high-energy heavy ion collisions shows unique properties such as large
535 radial collective flow, jet quenching in the created matter. Measurement of identified hadron
536 p_T spectra is one of the ways to extract information on hadron production mechanisms and
537 final state interactions between produced hadrons. Especially, we focus on (anti)proton p_T
538 spectra at low to high p_T (<6 GeV/c) to study the beam energy and system size dependences
539 of radial flow strength with the available various data sets; p+p, Cu+Cu, Au+Au at $\sqrt{s_{NN}}$
540 = 62.4 and 200 GeV. From a different point of view, the so-called intermediate p_T region
541 (2-5 GeV/c) is considered to have relatively large soft contribution affected by radial flow in
542 addition to hard jet component. This may be part of the origin of the baryon enhancement
543 in this p_T region.

544
545 We will present $\pi/K/p$ p_T spectra measured at PHENIX. Several model fits including hydro-
546 inspired one are used to extract freeze-out properties for comparison between different colli-
547 sions systems. we also discuss relative contributions of radial flow at partonic and hadronic
548 phases with complementary elliptic flow measurement.

Stopping effects in U+U collisions with a beam energy of 520 MeV/nucleon

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549 The Cooling Storage Ring (CSR) External Target Facility (ETF) at Lanzhou, China[1],
550 delivering the uranium beam up to 520 MeV/nucleon is expected to make significant contri-
551 bution and provide a good opportunity to systematically study the the nuclear equation of
552 state (EOS)at the high net-baryon density region of nuclear matter phase diagram[2].

553 Uranium is the largest deformed stable nucleus, and has approximately an ellipsoid shape
554 with the long and short semi-axis given by $R_l = R_0(1 + 2\delta/3)$ and $R_s = R_0(1 - \delta/3)$,
555 respectively, where $R_0 = 7$ fm is the effective spherical radius and $\delta = 0.27$ is the deforma-
556 tion parameter. we consider two extreme orientations: the so-called tip-tip and body-body
557 patterns with the long and short axes of two nuclei are aligned to the beam direction, respec-
558 tively. With the two extreme collision orientations, some novel stopping effects which are
559 believed responsible for some significant experimental observables, such as particle produc-
560 tion, collective motion as well as attainable central densities, can be obtained. The various
561 collision patterns could be selected from random orientation collisions in certain purity and
562 efficiency by cutting On-Line hardware trigger such as mid-rapidity charge multiplicity and
563 forward neutron multiplicity, along with the Off-Line trigger of stopping power ratio.

564 A Relativistic Transport Model (ART1.0)[3] is applied to simulate the stopping effects
 565 in tip-tip and body-body U+U collisions, at a beam kinetic energy of 520 MeV/nucleon.
 566 Our simulation results have demonstrated that both central collisions of the two extreme
 567 orientations can achieve full stopping, and also form a bulk of hot, dense nuclear matter
 568 with a sufficiently large volume and long duration, due to the largely deformed uranium
 569 nuclei. The nucleon sideward flow in the tip-tip collisions is nearly 3 times larger than that
 570 in body-body ones at normalized impact parameter $b/b_{max} < 0.5$, and that the body-body
 571 central collisions have a largest negative nucleon elliptic flow $v_2 = -12\%$ in contrast to zero
 572 in tip-tip ones. The Large nucleon sideward flow in tip-tip collisions and the significant
 573 negative nucleon elliptic flow in body-body central collisions can provide a sensitive probe
 574 to extract nuclear EoS information. Thus the extreme circumstance and the novel collective
 575 flow in tip-tip and body-body collisions can provide a good condition and sensitive probe
 576 to study the nuclear EoS, respectively. More experimental observables of U+U collision
 577 dynamics should be further studied, due to the geometric effects state (EoS).

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PHENIX measurements of $\pi^+\pi^+$ and $\pi^-\pi^-$ 4-momentum correlation functions in Au+Au collisions at $\sqrt{s_{NN}} = 200\text{GeV}$

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582 4-momentum correlation functions of pion pairs in central $Au + Au$ collisions provide an
 583 important avenue for study of the spacetime extent of the QCD matter produced at RHIC. In
 584 a recent high statistics experimental run, the PHENIX collaboration has performed detailed
 585 measurements with improved resolution of the reaction plane. This allows for a detailed
 586 correlation analysis of mid-rapidity pions, $0.20 < p_t < 0.40$, GeV/C , in the 10% most
 587 central collisions with respect to the reaction plane. These measurements are essential for
 588 understanding the reaction plane dependence of the 3D source functions and gaining further
 589 insight into the non-Gaussian tails recently observed in the pion pair transverse momentum
 590 and beam direction. The latest results from our ongoing analysis will be presented.

Systematics study of Photon production at forward rapidity in STAR

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592 We present the measurement of multiplicity and pseudorapidity distributions of inclusive
 593 photons production at forward rapidity (η from -2.3 to -3.8) for Au+Au collisions at
 594 200 GeV and Cu+Cu collisions at 200 and 62.4 GeV. Photons are measured using the
 595 Photon Multiplicity Detector (PMD) in STAR experiment at RHIC. The distributions of
 596 (η_{beam}), where η_{beam} is beam rapidity will be presented. Pseudorapidity distributions of
 597 photons at forward rapidities, when observed in the frame of one of the colliding particles
 598 show the signature of energy independent scaling at forward rapidities known as Limiting
 599 Fragmentation. Measurement of inclusive photons reflect the distribution for pi0s as most
 600 of the photons come from pi0 decay. Gluon saturation at initial conditions of the collisions
 601 is one of the possible explanations for this scaling. Study of system-size and incident energy
 602 dependence of this phenomenon will help to understand the mechanisms better. The system-
 603 size and energy dependence of the widths of scaling in pseudorapity and comparison with
 604 model predictions will be presented.

Investigation of high p_t events and shear viscosity in Nucleus-Nucleus collisions using the Hijing event generator

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605 In the recent years lot of interest has been observed in the Nucleus-Nucleus collisions at
 606 RHIC energies in phenomena related to high p_t physics[1] that is the suppression of high p_t
 607 particles and disappearance of back-to-back jets compared to the scaling with the number
 608 of binary nucleon- nucleon collisions indicating that a nearly perfect liquid is produced in
 609 these collision. Results on self shadowing of high p_t events will be presented using hadron
 610 multiplicity associated to high p_t events and unbiased events in Nucleus-Nucleus collisions[2]
 611 obtained from the hijing event generator. The results on shear viscosity using transverse
 612 momentum correlations[3] will also be presented.

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ELLIPTIC FLOW METHODS IN ALICE

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619 Anisotropic flow is recognized as one of the main observables providing information on the
620 early stage of a heavy-ion collision. Anisotropic flow is characterized by a Fourier expansion
621 of the azimuthal particle distribution. The first two harmonics, v_1 and v_2 , are called directed
622 and elliptic flow respectively. Initial measurements of elliptic flow at RHIC have shown that
623 the system created in the collision to first order can be described as a thermalized fluid which
624 expands hydrodynamically.

625 It is expected that at the LHC, due to high initial energy densities, all elliptic flow is generated
626 before hadronization takes place; making v_2 a probe of the equation of state of the Quark
627 Gluon Plasma. A detailed measurement of v_2 , and higher harmonics, is possible at the LHC
628 because of the expected increase in multiplicity and in the magnitude elliptic flow itself.

629 For detailed measurements of anisotropic flow there are several methods available; the event-
630 plane, cumulants and Lee-Yang zeroes methods. These methods all have different sensitiv-
631 ities to non-flow and flow fluctuations. Combined they give a better estimate of the flow
632 signal. We present flow analysis using these methods on data simulated within the ALICE
633 framework.

634 Until recently only the event-plane method was used in ALICE for the second order reaction
635 plane estimation. We present here a new method to extract the reaction plane, based on the
636 Lee-Yang zeros technique. This event plane does not need to be corrected for autocorrelations
637 and is not influenced by non-flow. Therefore it provides a better estimate of the reaction
638 plane. The method is applied on simulated data to test its range of applicability.

Quantitatively Constraining Theory Input from Experimental Data with Statistical and Systematic Errors

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639 The PHENIX experiment has measured suppression of semi-inclusive single π^0 in Au+Au
640 collisions at $\sqrt{s_N N}=200$ GeV. The present understanding of this suppression is in terms
641 of parton energy loss in a dense color-charge medium. We present a quantitative method
642 for constraining parton energy-loss model parameters from our experimental data that have
643 statistical point-to-point uncorrelated errors and correlated systematic errors. We detail
644 this methodology and apply it to various data sets and models to extract quantities such as
645 medium transport coefficients.

$D_{sJ}(2317)$ meson production at RHIC

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646 Production of $D_{sJ}(2317)$ mesons in relativistic heavy ion collisions at RHIC is studied [1]. Us-
647 ing the quark coalescence model, we first determine the initial number of $D_{sJ}(2317)$ mesons
648 produced during hadronization of created quark-gluon plasma. The predicted $D_{sJ}(2317)$
649 abundance depends sensitively on the quark structure of the $D_{sJ}(2317)$ meson. An order-of-
650 magnitude larger yield is obtained for a conventional two-quark than for an exotic four-quark
651 $D_{sJ}(2317)$ meson. To include the hadronic effect on the $D_{sJ}(2317)$ meson yield, we have
652 evaluated the absorption cross sections of the $D_{sJ}(2317)$ meson by pion, rho, anti-kaon,
653 and vector anti-kaon in a phenomenological hadronic model. Taking into consideration the
654 absorption and production of $D_{sJ}(2317)$ mesons during the hadronic stage of heavy ion col-
655 lisions via a kinetic model, we find that the final yield of $D_{sJ}(2317)$ mesons remains sensitive
656 to its initial number produced from the quark-gluon plasma, providing thus the possibility
657 of studying the quark structure of the $D_{sJ}(2317)$ meson and its production mechanism in
658 relativistic heavy ion collisions.

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Rapidity Dependence of Coalescence in Au+Au Collisions at $\sqrt{s_{NN}} = 200$ GeV.

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The BRAHMS Experiment at RHIC measures charged particles at a wide range of deuteron rapidities, $y \sim 0 - 3$, offering a unique opportunity to study the rapidity dependency of various physical phenomena.

In this presentation deuteron coalescence is used to study the rapidity dependence of the source size. As the fireball cools down and hadronizes it becomes possible for protons and neutrons to coalesce into deuterons. The coalescence parameter,

$$B_2 = \frac{E_d \frac{d^3 N_d}{dp_d^3}}{(E_p \frac{d^3 N_p}{dp_p^3})^2}$$

662 can be measured and connected via model calculations, to the volume of the fireball. Thus
663 measurements of B_2 can yield information about the density profile of the emitting source
664 at the time of hadronic freeze-out.

665 We present here the rapidity, as well as centrality, dependence of B_2 in $\sqrt{s_{NN}} = 200$ GeV
666 Au+Au collisions. The obtained results are compared with similar published measurements
667 from other RHIC experiments. Furthermore an estimate of the source size radius is given.

**Measurement of ω mesons via radiative decay mode
in $\sqrt{s_{NN}}=200\text{GeV}$ Au+Au collisions at RHIC-PHENIX**

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668 Measuring low mass vector mesons in an extreme condition made by relativistic heavy-ion
669 collisions is one of the intriguing study as a quest for the QCD phase transition to Quark-
670 Gluon-Plasma. Properties such as mass and width are expected to be modified at low p_T as
671 a probe of the chiral symmetry restoration. Furthermore, hadron suppression due to the jet
672 quenching should be seen in their spectra.

673 We report on ω mesons' production via radiative decay mode ($\omega \rightarrow \pi^0\gamma, \pi^0 \rightarrow 2\gamma$) in $\sqrt{s_{NN}}$
674 GeV Au+Au collisions at RHIC-PHENIX. The main challenge of this analysis is to handle the
675 huge combinatorial background inevitable for reconstructing particles from the three-body
676 decay mode. The emphasis is on the feasibility study using both simulation and real data to
677 search the best parameters that improve S/\sqrt{B} . For example, momentum and energy range
678 of π^0 and γ is optimized to reconstruct ω . To estimate the remaining background, we consider
679 three background sources; correlated/uncorrelated background and K_s^0 contribution. Then
680 each amount of background is determined by fitting and the total background is subtracted
681 successfully. The result suggests that ω meson production is suppressed in central Au+Au
682 collisions at high p_T region.

ALICE Transition Radiation Detector

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683 ALICE (A Large Ion Collider Experiment) has been designed as a dedicated heavy-ion
684 experiment and its primary aim is the study of nuclear collisions at LHC. In order to establish
685 and analyse the existence of the quark gluon plasma (QGP), a number of observables have
686 to be studied with ALICE in a systematic and comprehensive ways. Measurements of heavy
687 quarkonia states, such as J/Ψ , Ψ' , Υ , Υ' , and high p_T charged particles carry important
688 information to characterize the QGP. Transition Radiation Detector (TRD) was developed
689 to improve the electron identification capability and the tracking resolution in the central
690 rapidity region of the ALICE coverage. Using the TRD, reasonable J/Ψ and Υ measurements
691 will be realized[1].

692 The TRD surrounds the TPC (Time Projection Chamber) in the central barrel of ALICE.
693 It consists of 540 drift chambers with the drift direction perpendicular to the wire planes.
694 Xe,CO₂(15%) gas is used to have ionization energy deposit and absorption of X-rays from
695 transition radiation. The size of the drift chamber is $\sim 1.2 \times 1.4 \text{ m}^2$ and the total number of
696 readout channels is 1.18 million. Drift chambers are arranged in 18 large scale (7 m length)

697 modules (super-modules). One super-module contains 5 longitudinal stacks and 6 radial
698 layers of chambers each. With 6 layers, TRD achieves less than 2% pion efficiency with
699 keeping 90% electron efficiency.

700 The development of custom on-detector front end electronics with ASIC, multi-chip module,
701 and chip-to-chip bonding technologies has been finished and integrated with the detectors.
702 Two super-modules have been assembled and installed into a magnet in ALICE. The cosmic
703 run will be performed in December together with the other ALICE detectors and infrastruc-
704 tures. In addition, the test beam experiment is being performed at CERN PS using electron
705 or pion beam.

706 The construction technology of TRD super-module, electronics and control system, and the
707 basic performance of the super-module obtained in the test beam experiment and cosmic
708 run are described in this poster.

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Continuum Thermodynamics of $SU(3)$ Gauge Theory in $2 + 1$ vs $3 + 1$ Dimensions

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712 We present the thermodynamical variables in a controlled continuum extrapolation of finite
713 temperature $SU(3)$ lattice gauge theory in $2 + 1$ dimensions. We have investigated in detail
714 systematic errors that come from finite size and finite lattice spacing. The $2 + 1$ dimensional
715 theory is easy to simulate on the lattice compared with the $3 + 1$ dimensional theory, and,
716 of course, QCD with fermions. It is, however, in many aspects similar to pure gauge theory
717 in $3 + 1$ dimensions, it has linear confinement and a deconfining phase transition at finite
718 temperature. We compare our results with resummed perturbation theory. We also give a
719 simple analytic interpolation of our data in the range $T_c \leq T \leq 10T_c$, which can be used
720 to compare the results with phenomenological models. Finally we discuss the implications
721 of our analysis for the thermodynamics of a strongly interacting quark gluon plasma in $3+1$
722 dimensions.

Muon measurements with the CBM Experiment

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723 The Compressed Baryonic Matter (CBM) experiment at the Future Facility for Antiproton
724 and Ion Reaseach(FAIR) is aiming at the investigation of strongly interacting matter at
725 high baryon densities. The research program of CBM comprises of the exploration of the
726 equation of state of dense hadronic matter, the search for the deconifinement first order phase
727 transition which is expected to occur at high net-baryon densities and the study of in-meium
728 modifications of hadrons which might signal the onset of chiral symmetry restoration.

729 The key observables include vector mesons decaying into lepton pairs. In particular J/ψ and
730 ψ' mesons are predicted to be sensitive probes of the deconfined phase. The rare probes will
731 be measured for the first time at FAIR energies. Due to the high reaction rate(up to $10MHz$)
732 envisaged for the CBM experiment, the recorded charmonium yields will be comparable to
733 measurements at LHC.

734 Results of feasibility studies will be presented for the identification of charmonium via its
735 dimuon decay in heavy-ion-collision at FAIR energies. Moreover detector requirments and
736 possible technological solutions will be discussed. Initial R&D on detector development will
737 be presented.

Pion form factor in asymmetric nuclear matter

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738 It is well known that the pion form factor ($F(Q^2)$) plays an important role to describe the
739 photon-hadron interactions in vector meson dominance (VMD) model. $F(Q^2)$ is proportional
740 to the rho spectral function with a peak at the ρ pole mass and a kink near the ω pole
741 indicating the mixing of ρ - ω meson in free space due to explicit symmetry breaking driven
742 by u - d quark mass difference. There exists several models which calculate mixing amplitude
743 of these mesons and their momentum dependence.

744 In the present work we study the medium modification of pion form factor in asymmetric
745 nuclear matter within the framework of Quantum Hadrodynamical model. It is shown that
746 both the shape and the pole position of the pion form factor in dense asymmetric nuclear
747 matter is different from its vacuum counterpart with ρ - ω mixing. This is due to the density
748 and asymmetry dependent ρ - ω mixing which could even dominate over its vacuum coun-
749 terpart in matter. Effect of the in-medium pion factor on experimental observables *e.g.*,

First Result of Net-Charge Jet-like Correlations from STAR
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751 Two-particle jet-like correlations of charged hadrons with a high p_T trigger particle are
 752 strongly modified in central Au+Au collisions at RHIC, lending strong support for jet
 753 quenching and partonic energy loss. At intermediate p_T the away-side correlation exhibits a
 754 double-peak structure which prompted many theoretical investigations of jet-medium inter-
 755 actions. In the same intermediate p_T range the inclusive baryon to meson ratio was found
 756 to be strongly enhanced in central Au+Au collisions and is significantly higher than that
 757 from jet fragmentation in elementary collisions. Studying jet-like correlations separately for
 758 baryons and mesons, therefore, will shed more light on the physics of jet-medium interaction.
 759 To this end, we analyze jet-like correlations with associated positive and negative charged
 760 particles separately. The net-charge (difference between positive and negative charged parti-
 761 cles) may be dominated by net-protons, therefore is primarily sensitive to the baryon content,
 762 while the total charge contains all charged particles and has a very different baryon to meson
 763 composition than the net-charge. We present the first results from STAR on net-charge jet-
 764 like correlations in Au+Au collisions as a function of centrality and compare to total-charge
 765 results as well as results from reference d+Au data. Associated p_T spectra of net-charge
 766 yields on both the near and away side are obtained, and are compared to those of total
 767 charge correlated yields as well as inclusive yields. Ratio of net-charge to total charge cor-
 768 related yields is presented as a function of associated p_T and centrality, and is compared to
 769 the corresponding ratio from inclusive measurements which reflects the medium. The results
 770 will be discussed in the context of jet-medium interactions at RHIC.

Simulation of first order deconfinement transition in relativistic heavy-ion collisions

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771 We have carried out numerical simulation of the first order confinement-deconfinement phase
 772 transition to study the formation and evolution of Z(3) domain walls in relativistic Heavy
 773 Ion Collision. We use the effective potential for the Polyakov loop order parameter for this
 774 transition, proposed by Pisarski. The transition in this model is weakly first order, and

775 we simulate it by first determining the profile of the critical bubble using the bounce tech-
776 nique and then simulating the phase transition by random nucleation of these bubbles. The
777 nucleation probability from the initial confined vacuum to the deconfined state in the con-
778 fining background is used from the literature to determine the number of bubbles nucleated
779 per unit time in the overlap region of the colliding nuclei. This initial field configuration
780 is evolved using a stabilized leapfrog algorithm of second-order accuracy. Due to Bjorken
781 scaling for the longitudinal expansion, early dynamics of bubble evolution is described using
782 simulation in 2+1 dimensions. As the bubbles expand and coalesce, $Z(3)$ walls and associ-
783 ated strings form due to random choice of $Z(3)$ vacua inside different bubbles. We study
784 the coarsening of the domain wall network and any possibility of scaling solution within the
785 short time available for the system before the hadronization stage when the walls melt away.
786 We also estimate the energy density fluctuations resulting from the decay of walls when
787 temperature drops below T_c . We use this simulation to make a detailed study of observable
788 consequences of the presence of these walls in terms of enhancement of high P_T hadrons,
789 and multiplicity fluctuations resulting from the decay of wall network, as well as presence of
790 novel jet structures.

Transverse Energy Measurement in Au+Au collisions by the STAR Experiment.

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791 Transverse energy (E_T) has been measured with both of its components, namely hadronic
792 (E_T^{had}) and electromagnetic (E_T^{em}) at mid-rapidity, for 62.4 GeV and 200 GeV (Run-IV)
793 Au+Au collisions by the STAR experiment. In the common phase space of TPC and BEMC
794 ($0 < \eta < 1$ and full ϕ), we have obtained E_T^{had} from the TPC reconstructed tracks and
795 E_T^{em} from the BEMC tower hits after correcting for the hadronic contaminations in the
796 calorimeter. We present the minimum-bias spectra of E_T and its components for 62.4 GeV
797 and 200 GeV Au+Au collisions. The centrality behavior of $\langle dE_T/d\eta \rangle / (0.5N_{part})$ and
798 the excitation function of $\langle dE_T/dy \rangle / (0.5N_{part})$ has been compared with the final state
799 gluon saturation model (EKRT) along with similar measurements from SPS to RHIC. The
800 centrality behavior of the ratio $\langle dE_T/d\eta \rangle / \langle dN_{ch}/d\eta \rangle \equiv E_T/N_{ch}$ has been compared
801 to the corresponding data at other energies. The most striking feature is the observation of
802 a nearly constant value of $E_T/N_{ch} \sim 0.8$ GeV from AGS, SPS to RHIC. It has been shown
803 by the statistical-thermal model that this behavior is related to the freeze-out criteria and
804 is consistent with the single freeze-out model [1]. It is found that the increase in E_T produc-
805 tion with energy, comes mostly from the increase in particle multiplicity. The E_T and N_{ch}
806 production has been observed to follow a logarithmic behavior. Based on the boost-invariant
807 Bjorken hydrodynamic model, the estimation of the initial energy density of the produced
808 fireball, when compared with the lattice QCD calculations, suggests that a deconfined phase
809 has been created at RHIC for the central Au+Au collisions. Taking similar colliding species

810 i.e. Au+Au, the $\epsilon_{Bj} \cdot \tau$ has been predicted for LHC, based on the measurements at RHIC.
 811 We have also made predictions for $(dE_T/d\eta)/(0.5 N_{part})$ and $(dN_{ch}/d\eta)/(0.5 N_{part})$ for the
 812 LHC energy. The ratio E_T^{em}/E_T , for top 5% central events is consistent with a final state
 813 dominated by mesons and has been found to increase slowly with $\sqrt{s_{NN}}$. E_T^{em}/E_T has been
 814 observed to be independent of centrality. Furthermore, we have studied the event-by event
 815 fluctuations in E_T and in the ratio of it's components using σ/μ as the fluctuation observ-
 816 able. The fluctuations in $E_T, E_T^{em}/E_T$ have been observed to be 8.4% and 9.7% respectively
 817 for top 5% central events. The centrality behavior of E_T -fluctuation has also been studied.

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Track Matching Study from TPC to PHOS in ALICE Experiment

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820 The Photon Spectrometer (PHOS) designed by ALICE is dedicated to measuring photon spectrum
 821 precisely with a large transverse momentum range from 0.5-100GeV/c. We present a study within
 822 the frame of software-AliRoot for matching charged tracks reconstructed from the ALICE tracking
 823 detectors with the clusters measured in the PHOS. Matching efficiency and contamination due to
 824 wrong matches have been deduced for charged pions, muons and protons. For electrons, lower
 825 efficiency and more contamination are discovered. And further the study shows that with less
 826 additional material in front of PHOS will be much in favor of the photon detection and thermal
 827 photon analysis of the middle transverse momentum domain.

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Measurement of High p_T Flow and Jet Correlation in Au+Au Collisions at $\sqrt{s_{NN}} = 200 GeV$ from PHENIX

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836 Single hadron yields at high transverse momentum p_T are observed to be strongly suppressed in
 837 central Au+Au collisions. This suppression is thought to be the result of jet quenching, and the
 838 observed high p_T hadrons are thought to be dominated by fragmentation of partons emitted from
 839 the surface of the overlap region. Although extremely simplistic, this picture would lead to a small

840 but finite v_2 and a relatively unmodified near-side jet shape for these high p_T hadrons. A combined
841 study of the v_2 and near-side jet correlations of leading hadrons can improve our understanding of
842 the particle production mechanism and geometrical bias in this p_T region. PHENIX experiment
843 installed a new reaction-plane detector and collected 5.5 billion Au+Au events in RUN7. The
844 significantly improved ability to determine the reaction-plane and large statistics allow a more
845 detailed study of flow and jet correlation. Results on the elliptic flow and jet correlation for high
846 p_T hadrons are discussed.

Analysis of Low Energy Pion Spectra

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847 The transverse mass spectra and the rapidity distributions of π^+ and π^- in Au-Au collisions at 2,
848 4, 6, and 8GeV·A by E895 collaboration are fitted using an elliptically expanding fireball model
849 with the contribution from the resonance decays and the final state Coulomb interaction. The
850 ratio of the total number of produced π^- and π^+ is used to fit the data. The resulting freeze-out
851 temperature is rather low($T_f < 60\text{MeV}$) with large transverse flow and thus resonance contribution
852 is very small. The difference in the shape of m_t spectra of the oppositely charged pions are found
853 to be due to the Coulomb interaction of the pions.

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857

FLOW FLUCTUATIONS AND MEASUREMENT OF ANISOTROPY COEFFICIENTS

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858 The observation of flow in relativistic heavy ion collisions has led to different analysis techniques to
859 determine flow from experimental data. The standard method determines the coefficients of fourier
860 expansion of azimuthal distribution by obtaining the average cosines of the particle angles with
861 respect to the estimated reaction plane. This method measures essentially two particle correlations,
862 and the coefficients include contributions from non-flow correlations. The methods of obtaining the
863 higher order cumulants and that of finding the Lee-Yang zeroes determine many particle correlations

864 due to collective effects.

865 Considering the recent interest in fluctuations in flow we propose to revisit the standard method of
866 analysis. The different implementations of the standard method of analysis yield consistent results.
867 The method purports to determine the mean value of flow, $\langle v_n \rangle$, for any given sample of events,
868 where as it determines $\sqrt{\langle v_n^2 \rangle}$. The two values agree with each other if the values of flow for all
869 events are distributed about a mean with a *small* variance.

870 We shall argue that the dependence of flow on centrality, rapidity and particle identity introduces
871 an additional uncertainty in the p_T dependence of v_2 , manifesting in different values of flow for
872 each event. This mimics fluctuations in flow, and contributes to the variance of the distribution
873 over and above the contribution due to impact parameter and eccentricity fluctuations. The effect
874 of this uncertainty in the p_T dependence will be estimated on the final values of flow determined
875 using the standard method of analysis.

876 The standard method often decreases the contribution of non-flow correlations by determining the
877 flow of a given set of particles using the estimated reaction plane from another set of particles. We
878 shall also show that fluctuations in the flow of particles used to estimate the reaction plane also
879 contribute to similar behaviour, namely, decreasing the measured values of flow for the first set of
880 particles.

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DILEPTON FROM QGP AND QUARK-HADRON PHASE TRANSITION

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886 A model of statistical quark-gluon plasma formation is considered. We look the dilepton pro-
887 duction at the transition temperature which is obtained around (150 – 170) *MeV*, from this
888 system with the quark mass depending on the coupling parameter used in R. Ramanathan et
889 al(*PRC70*, 2004, 027903). The rate of dilepton production is shown for the invariant mass M as
890 well as transverse momentum P at the particular value of $E = 2.0, 3.0$ *GeV*. It is found that there
891 is significant production of leptons in this process.

Measurement of π^0 and η mesons with PHENIX in sqrt(s)=200GeV Au+Au collisions at RHIC

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892 The π^0 meson has been a crucial probe for observing jet quenching in ultrarelativistic heavy-ion
893 collisions at RHIC. With the data from the 2004 RHIC run, a first systematic comparison between
894 the precise measurement of the nuclear modification factor (R_{AA}) and theoretical calculations al-
895 lowed constraining model parameters such as the initial gluon density dN^g/dy and the transport

896 coefficient \hat{q} . Measurements of the η meson production in the same collisions have also shed light
897 on a possible dependence of the suppression on the particle species.

898

899 In the 2007 Au+Au run of RHIC, PHENIX has gathered more than 3 times the events than in
900 2004, allowing an extension of the particle spectra up to higher transverse momenta. PHENIX has
901 also added new detectors to its setup that allow higher precision e.g. in the measurement of the
902 reaction plane. The latest result on the measurement of the π^0 and η mesons as well as the status
903 of the Run7 data analysis will be shown and discussed.

D⁰ production in proton proton collision at $\sqrt{s} = 200$ GeV

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904 The charm quarks are a rich probe of QCD matter produced in heavy ion collisions. To understand
905 how the energy loss mechanism and the flow strength depend on quark mass, one has to study
906 charm quark production. A systematic study of charm yield and spectra in pp, pA and AA is
907 a prerequisite to understand the properties of the strongly interacting matter. In this work, we
908 report on the D⁰ reconstruction in the K π channel in pp collisions at $\sqrt{s} = 200$ GeV, measured
909 by PHENIX collaboration. D⁰-mesons are produced in pairs. While one of them may decay to
910 K π the other may decay to electron channel. The electron decay provides a method of tagging
911 the D⁰ mesons and helps reducing the hadronic background from pp. The K and π identification
912 is optimized with a combination of time of flights in TOF and EMCAL subsystems. Mixed event
913 technique is used to subtract the background. The simulations are carried out to study the effect
914 of electron tagging on signal to background ratio and to tune the hadron identification. A status
915 of this analysis will be presented.

Azimuthal Correlations of Electrons from Heavy Quark Decay with Hadrons in Au+Au Collisions at $\sqrt{s_{NN}}=200$ GeV in PHENIX

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916 The observed large suppression and elliptic flow of electrons from heavy quark decay has led to
917 a rethinking of the conventional radiative energy loss picture that well describes light quark en-
918 ergy loss. However, substantial ambiguities exist because of uncertainties in the charm/bottom
919 ratio and unknown medium effects on the fragmentation and hadronization of heavy quarks. Two
920 particle hadron correlations have been used extensively to obtain additional information about
921 hadronization, fragmentation and possible medium response in heavy ion collisions. The PHENIX
922 experiment is well suited to electron measurements because of it's excellent electron identification
923 capabilities. We apply techniques developed for direct photon-hadron correlations to statistically
924 subtract correlations from photonic sources from the inclusive electron hadron correlations and we
925 present the current status of two particle correlations between electrons from heavy quark decay
926 and charged hadrons in Au+Au collisions in the PHENIX experiment.

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929 Understanding the J/ψ suppression and possible recombination mechanisms at RHIC is one of
930 the outstanding challenges for theorists and experimentalists. Recent results provided by PHENIX
931 showed a stronger suppression at forward rapidities, while at mid-rapidity the suppression is similar
932 to lower energy collision experiments. A large sample of Au+Au collisions at $\sqrt{S_{NN}}=200$ GeV was
933 collected during 2007 data taking run with PHENIX detector at RHIC. Using this sample, J/ψ s
934 were identified in the di-electron decay channel. We studied the J/ψ nuclear modification factor
935 and the first determination of its v_2 elliptic flow parameter at mid-rapidity.

K*(892) RESONANCE PRODUCTION IN PB+PB COLLISIONS AT CERN SPS ENERGIES

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976 Production of the $K^*(892)$ resonance in the $K^+\pi^-$ decay channel was studied by the NA49 exper-
977 iment in central Pb+Pb collisions at 158A GeV beam energy. The signals of the $K^*(892)$ mesons
978 are extracted from the invariant mass spectra of $K^+\pi^-$ pairs using particle identification by dE/dx
979 measurements in the NA49 TPCs. The mixed event background subtraction technique was opti-
980 mised for dealing with the overwhelming combinatorial background. The mass and width of the
981 resulting mass peak are consistent with the PDG table values. Preliminary $K^*(892)$ yields are
982 presented as a function of rapidity and transverse momentum.

Jet Quenching in Relativistic Heavy Ion Collisions

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983 Quantum chromo dynamics (QCD) calculations on space-time lattice predict a transition from
984 confined state in hadrons to a deconfined state of quarks and gluons above a temperature of 170
985 MeV or an energy density above $1\text{GeV}/\text{fm}^3$. In the relativistic heavy ion collisions at RHIC at
986 BNL, where the initial energy density is around $5\text{GeV}/\text{fm}^3$, such a transition to a deconfined state
987 of matter is formed as shown by experimental measurements of several observables [1,2]. It is
988 currently believed that this deconfined state consists of a strongly interacting Quark Gluon Plasma
989 (sQGP), behaving nearly like a perfect liquid, see [1,2]. Among many observables, jet-quenching
990 phenomenon is so far an undisputed and important signal for a hot dense medium formed in
991 relativistic heavy ion collisions. We are calculating the energy loss of partons while traversing
992 the QGP medium due to the gluon bremsstrahlung processes. The coherent radiation processes
993 involve multiple scatterings of the partons in the QGP medium during the gluon formation time.
994 This leads to the interference effects known as Landau-Pomeranchuk-Migdal effect (LPM) [3]. For
995 this purpose, we follow the Arnold Moore Yaffe (AMY) formalism discussed in detail in [3,4,5].
996 The bremsstrahlung integral equations determine the gluon emission amplitude given by $\mathbf{F}_s(\mathbf{h},\mathbf{p},\mathbf{k})$
997 function. In terms of this $\mathbf{F}_s(\mathbf{h},\mathbf{p},\mathbf{k})$ function, the emission rate is determined as given in Eq. 1
998 (for details see Eq.5 given in Ref.4). In our recent works, we used the iterations method and the

999 variational method for solving similar integral equations for virtual photon emission from QGP [6].
 1000 For the case of gluon emission, we are using the iterations method to solve AMY integral equation
 1001 to obtain $\mathbf{F}_s(\mathbf{h},p,k)$ function. The work is in progress and in the conference; we will present results
 1002 for $\mathbf{F}_s(\mathbf{h},p,k)$ distributions.

$$\frac{d\Gamma_g^{LPM}}{d^3\mathbf{k}} = \frac{\alpha_s}{4\pi^2 k^2} \sum_s N_s d_s C_s \int_{-\infty}^{+\infty} \frac{dp}{2\pi} \int \frac{d^2\mathbf{h}}{4\pi^2} n_s(p+k)[1 \mp n_s(p)] \otimes [1 + n_b(k)] \frac{1}{k^3} |\Gamma_{p \rightarrow p+k}^s|^2 2\mathbf{h} \cdot \Re \mathbf{F}_s(\mathbf{h}; p, k) \quad (1)$$

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PHENIX measurements of particle species dependent flow correlations in Au+Au and Cu+Cu collisions at $\sqrt{s_{NN}}=200$ GeV

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Flow correlation measurements constitute a central pillar in the body of evidence in support of the creation of the quark gluon plasma (QGP) in RHIC collisions. Such momentum anisotropy is now routinely characterized at mid-rapidity, by the even order Fourier coefficients,

$$v_n = \left\langle e^{in(\phi_p - \Phi_{RP})} \right\rangle, n = 2, 4, \dots, \quad (2)$$

1012 where ϕ_p is the azimuthal emission angle of a particle, Φ_{RP} is the azimuth of the reaction plane
 1013 and the brackets denote statistical averaging over particles and events.

1014 At RHIC energies, momentum anisotropy is believed to result from frequent particle interactions
 1015 which drive pressure gradients in an ‘‘almond-shaped’’ collision zone. However, the question of
 1016 whether or not particle collisions are frequent enough to achieve full local equilibrium is still under
 1017 debate. Similarly, a consensus on the detailed dynamical evolution of the QGP has not been
 1018 reached.

1019 In recent experiments, the PHENIX collaboration has made detailed measurements of flow correla-
 1020 tions and their fluctuation, for several particle species produced in Au+Au and Cu+Cu collisions
 1021 at $\sqrt{s_{NN}}=200$ GeV, with an eye toward resolving these issues.

1022 We will present a detailed set of harmonic flow correlation results and discuss their implication for
 1023 local thermalization and the dynamical evolution of the QGP, as reflected by its equation of state
 1024 and transport properties.

Probing the Quark-Gluon Phase Transition using Energy and System-Size Dependence of Long-Range Multiplicity Correlations in Heavy Ion Collisions from the STAR Experiment

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1025 Long-range forward-backward multiplicity correlations have been measured in heavy ion collisions at
1026 RHIC. Results for short and long-range multiplicity correlations (forward-backward) are presented
1027 for several systems (Au+Au, Cu+Cu, and pp) and energies (e.g. $\sqrt{s_{NN}} = 200$ and 62.4 GeV).
1028 A strong, long-range correlation is seen for central heavy ion collisions at $\sqrt{s_{NN}} = 200$ GeV that
1029 vanishes in semi-peripheral events and *pp* collisions. There is no apparent scaling with the number
1030 of participants (N_{part}) involved in the collision. These correlations provide information about
1031 the longitudinal behavior of the system formed in heavy ion collisions. To access the transverse
1032 behavior, the clusters produced in the same heavy ion collisions have been characterized by a
1033 study of the energy and system size dependence of the percolation density parameter (ρ) [1]. The
1034 relationship between the long-range correlation and ρ has been explored to characterize the hadron-
1035 quark/gluon phase transition and rapid thermalization of the system.

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Properties of the ϕ meson at high temperatures and densities

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1038 We calculate for the first time the mass shift, the width broadening, and the spectral density of the
1039 ϕ meson in a hot bath of nucleons and pions using a general formalism relating self-energy to the
1040 forward scattering amplitude (FSA) [1,2]. In order to describe the low energy FSA, we use exper-
1041 imental data along with a background term. For the high energy FSA, a Regge parametrization
1042 is employed. We verify the resulting FSA using dispersion techniques. We find that the position
1043 of the peak of the spectral density is a little shifted from its vacuum position, but its width is
1044 considerably increased due to collisional broadening. Properties of the spectral density are also
1045 investigated in various physical limits [3]. Finally, we study the implications of our findings for
1046 measurements at SPS and at RHIC.

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Centrality dependence of charged kaon spectra from Au+Au collisions at 200 GeV

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1051 Strangeness enhancement is a useful signal to study Quark Gluon Plasma in relativistic heavy ion
1052 collisions. Kaons created in the collisions carry a large fraction of strangeness, which is not in the
1053 initial colliding nucleus. Measurements of kaon spectra over a wide centrality range will allow us to
1054 study the freeze-out dynamics. In particular at the intermediate transverse momentum, centrality
1055 dependence of kaon spectra will allow us to study the hadron formation mechanism via coalescence.
1056 In this talk, we present the charged kaon spectra at mid-rapidity from Au+Au collisions at 200 GeV
1057 measured by the STAR experiment at RHIC. With Time-of-Flight detector to identify kaon, the
1058 centrality dependence of kaon spectra is presented at $0.3 \text{ GeV}/c < p_T < 5 \text{ GeV}/c$. The freeze-out
1059 dynamics indicated by kaon spectra measurements are discussed. The comparison with 62.4 GeV
1060 Au+Au results are also discussed.

1061 **Full jet reconstruction in $p + p$ at $\sqrt{s} = 200 \text{ GeV}$ and in $\text{Cu} + \text{Cu}$ at $\sqrt{s_{NN}} = 200 \text{ GeV}$ with**
1062 **the PHENIX detector**

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1063 Acceptance-limited detectors, such as the PHENIX Experiment, are an especially challenging envi-
1064 ronment to perform full jet reconstruction. The heavy ion background further limits the usefulness
1065 of traditional jet reconstruction schemes when applied at RHIC energies. Yet the information de-
1066 rived from reconstructed jets could provide effective means to study in-medium energy loss and
1067 medium-induced modification of fragmentation.

1068 We introduce a new jet reconstruction algorithm based on filtering that is suitable for both
1069 acceptance-limited and large acceptance detectors, and is in part inspired by pQCD energy flow
1070 variables [1]. We will discuss the jet reconstruction performance in both PHENIX and full ac-
1071 ceptance detectors, and will present the status of first results from this general method for jet
1072 production in both hadronic and heavy ion collisions at RHIC.

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Study on the γ ray and neutron sensitivity of MRPC

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1076 The multi-gap resistive plate chamber (MRPC) has good time resolution (less than 100ps) and high
1077 detection efficiency (higher than 95) and will be used to construct a full barrel time-of-flight (TOF)
1078 detector for the STAR experiment at RHIC. In this paper we try to use MRPC technology to detect
1079 γ rays and fast neutrons. DPF neutron source and flash x-ray source were used to measure the γ
1080 ray and neutron sensitivity. The results show that MRPC has good response to x-ray and it can
1081 discriminate the γ ray peak and neutron peak of DPF source.

1082 In the mean time, we studied the method to improve the neutron sensitivity of MRPC. Thin
1083 polyethylene layer is placed on the surface of electrode glass, when neutron interact with polyethy-
1084 lene, proton is produced and it can ionize working gas. Electron and ions will produce avalanche
1085 by strong electric field and charges is induced on collect electrodes. This kind of neutron detector
1086 not only has high neutron- γ resolution capability, but also has good spatial resolution.

Measurement of direct photon via internal conversion in $\sqrt{s} = 200$ GeV p+p collisions at RHIC-PHENIX

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1087 Direct photons are unique probes to investigate properties of Quark Gluon Plasma (QGP) since
1088 photons are not influenced by the strong interaction. Photons are emitted at many different stages
1089 of the collisions, from the initial interaction to the final hadron-gas phase. However thermally emit-
1090 ted photons from QGP phase are considered to be the primary contributor in the low energy region.

1091
1092 A measurement of 'real' direct photon using an electromagnetic calorimeter is notoriously difficult
1093 in the low energy region due to large systematic errors. There are three main reasons why the
1094 errors can not be reduced:

- 1095 1. It is essentially very difficult to measure the photon energy with high resolution in the low
1096 energy region using an EMCal.
- 1097 2. The uncertainty in estimating hadron contamination becomes larger in this region.
- 1098 3. There is large irreducible π^0 background in the real photon spectrum.

1099 On the other hand, it is possible to measure 'virtual' direct photons via internal conversion with
1100 small systematic errors in the low energy region since the hadron contribution can be estimated
1101 using the Kroll-Wada formula.

1102
1103 It has been shown that this method of measuring 'virtual' photons can reduce the systematic errors
1104 of the photon yield in Au+Au collisions. Thus the same method is applicable to the direct photon
1105 yield in p+p collisions, which serves as a vital comparison to the direct photon yield in Au+Au
1106 collisions. The comparison of direct photon yield between in p+p and Au+Au collisions can be
1107 performed to discuss the production of thermal photon in QGP phase. In this work, we present
1108 the results of the measurement of the direct photon yield via internal conversion in p+p collisions
1109 at $\sqrt{s} = 200$ GeV.

1110 Understanding the behavior of QCD matter under different temperature and density conditions
1111 is a subject of intense experimental and theoretical work in nuclear physics. Recent results at
1112 RHIC suggest the formation of an ultradense form of matter [1]. However, in order to properly
1113 interpret these results, methods for analyzing background nonequilibrium processes need to be
1114 developed so that these processes may be separated from those that occur at equilibrium. For
1115 example, although the quark-gluon plasma (QGP) is not expected to be formed in p+p, p+A,
1116 or d+A collisions, these collisions result in effects which must also be ascertained separately from
1117 those that may be associated with the QGP [2]. These cold nuclear matter effects, as well as the
1118 modification of initial parton distribution functions, such as shadowing, gluon saturation, etc., need
1119 to be understood before firm conclusions can be drawn on the nature of the hot nuclear medium
1120 that has been produced at RHIC.

1121 An incoherent binary nucleon-nucleon collision model of AA collisions for simulating particle pro-
1122 duction in cold nuclear matter is presented, which with a simple phenomenological parameter, the
1123 mean nucleon energy loss fraction, yields results that match experimental data [3], as well as those
1124 of HIJING [1,4], rather well. The model's simulation program extracts particle production data
1125 from pp collisions that are simulated using the PYTHIA program. The nuclear geometry is de-
1126 scribed through the Glauber Model. The preliminary R_{AA} and R_{CP} results will be presented in
1127 this poster.

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Meson-baryon femtoscopy in AuAu collisions at 200GeV measured by STAR experiment

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1137 Correlations between non-identical particles at small relative velocity probe asymmetries in the
1138 average space-time emission points at freezeout [1]. Such asymmetries may arise from long-lived
1139 resonances, bulk collective effects, or differences in the freezeout scenario for the different particle
1140 species. STAR has extracted pion-proton correlation functions from a high-statistics dataset of

1141 Au+Au collisions at $\sqrt{s_{NN}}$ GeV. We present a femtoscopic analysis of this data for all combi-
1142 nations of pions and (anti-)protons, for collisions of different centrality. The measurements are
1143 compared with calculations of a simple Blast-wave model, in which asymmetries are driven only by
1144 collective flow, as well as with Therminator [2], which also accounts fully for resonance effects.

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Proton femtoscopy in STAR

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1149 Femtoscopy of two particles investigates the properties of matter produced in heavy-ion collisions.
1150 It allows to study the space-time characteristics of the baryon-emitting medium. The distinctions
1151 in the scattering cross sections, the production of resonances and collective motion can create
1152 differences between space-time properties of proton and anti-proton emission. We present the
1153 latest results on baryon-(anti-)baryon femtoscopy performed by STAR for Au+Au collisions at
1154 $\sqrt{s_{NN}} = 200$ AGeV and analyze the centrality dependence of such observables. Hydrodynamics
1155 and similar models predict no emission asymmetry between particles of same masses, while models
1156 including rescattering in the final stage (UrQMD) show a small asymmetry. Using our data we
1157 attempt to distinguish between the two scenarios, using proton - anti-proton femtoscopy as a direct
1158 probe for the existence of final state rescattering.

EOS of QGP using Screened Cornell Potential

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1159 Equation of State [EOS] of Quark Gluon Plasma [QGP] using Screened Cornell potential based on
1160 Mayer's cluster expansion is presented here[1-4]. The Cornell potential $V(r_{ij})$ [5] is defined here
1161 for gluon plasma as

$$V(r_{ij}) = \left(\frac{3\alpha_s}{r_{ij}} - \sigma r_{ij} \right) e^{-m_D r_{ij}}$$

and for quark-antiquark plasma

$$V(r_{ij}) = \left(\frac{4}{3} \frac{\alpha_s}{r_{ij}} - \sigma r_{ij} \right) e^{-m_D r_{ij}}$$

where α_s , σ and m_D are strong coupling constant, string tension and Debye mass (inverse Debye screening length) respectively. In the QGP, the Debye mass of a chromoelectric charge follows from the static limit of the longitudinal polarization tensor which in the high temperature limit is given by

$$m_D^2 = \alpha_s^2 T^2 \left(1 + \frac{n_f}{6} \right)$$

1162 where n_f is the number of light quark flavors in the QGP [6]. In our numerical analysis, we adjusted
 1163 the two parameters σ and α_s such that we get qualitatively good fit to the EOS of LGT (Lattice
 1164 Gauge Theory). For all the flavors at high temperatures EOS figures coincide. At $T \approx T_c$ our
 1165 results do not match with the lattice results. This may be due to perturbative effects.

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Systematics of cut-off effects in lattice calculations with chemical potential

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1173 Recent lattice calculations have attained accuracies of 10% or more. Two factors are respon- sible
 1174 for this success: The discovery of new, powerful algorithms such as the Rational Hybrid Monte
 1175 Carlo (RHMC) algorithm, and the use of improved actions. In particular, at high temperature
 1176 and non-vanishing chemical potential, where today's lattice calculations are forced to be performed
 1177 on rather coarse lattices, the use of improved actions is mandatory. While improvement schemes
 1178 for thermodynamic calculations have been used for quite some time, little is known about the sys-
 1179 tematics of these effects in the presence of a non-vanishing chemical potential as well as in chiral
 1180 fermion formulations. The latter eliminate another problem in today's lattice studies of QCD ther-
 1181 modynamics, the explicit chiral symmetry breaking inherent in the staggered and Wilson fermion
 1182 formulations. With the rapid im- provement in computing resources chirally invariant discretiza-
 1183 tion schemes also become more and more attractive to be used for thermodynamic calculations.
 1184 Recently it has been questioned whether a non-vanishing chemical potential may introduce addi-
 1185 tional cut-off errors in the chirally invariant overlap fermion formalism [1]. We will show here that
 1186 this is not the case and explicitly calculate the cut-off induced errors in the high temperature,
 1187 Stefan-Boltzmann limit of QCD [2]. We extend this study to the case of do- main wall fermions
 1188 as well as truncated perfect actions of Wilson type. In this poster, we present explicit results for

1189 the size of cut-off effects at finite temperature and non-vanishing chemical potential that arise in
1190 various discretization schemes used for the fermion sector of QCD in lattice calculations.

1191 **References**

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1195 duction in studies of QCD thermodynamics at non-zero temperature and chemical potential, in preparation

1196 **Multi-strange hadron production measured by the STAR detector in Cu+Cu** 1197 **collisions at $\sqrt{s_{NN}} = 62.4$ GeV**

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1198 The study of particle production in relativistic heavy ion collisions is an important tool to
1199 understand the fireball dynamics and the various process present in such physical system. The
1200 production of strange hadrons is of special interest given the particular value of the strange quark
1201 mass and the abundance of these hadrons in such collisions, allowing for a statistical
1202 interpretation of the results. The transverse momenta spectra and the total yields of strange
1203 hadrons have been measured by the STAR experiment for Au+Au collisions at center of mass
1204 energies of 62.4 GeV, 130 GeV and 200 GeV per nucleon pair. All these measurements allowed
1205 not only a systematic study of particle production dependence on collision energy but on system
1206 size as well due to the different collision centralities that can be probed. In 2005, STAR has
1207 measured Cu+Cu collisions at 62.4 and 200 GeV in order to enhance these systematic studies on
1208 the energy and the system size dependence of particle production. The results for 200 GeV were
1209 already presented in previous conferences. We will present details and preliminary results of the
1210 analysis of Ξ and Ω and their anti-particles production for Cu+Cu collisions at $\sqrt{s_{NN}} = 62.4$
1211 GeV. These particles are identified through the reconstruction of their weak decay topology. The
1212 data presented in this poster enriches the systematic study of energy and system size dependence
1213 of yields and ratios between the AGS experiments to the top RHIC energy measurements.

Strange Hadron Elliptic Flow v_2 from 200 GeV Cu+Cu Collisions

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1214 Elliptic flow allows us to probe early dynamics in high-energy nuclear collisions. Previous studies

1215 of elliptic flow, in particular its p_T dependence, of all charged particles and identified hadrons [1,2]
1216 from 200 GeV AuAu collisions at RHIC suggest that, the system has gone through early
1217 thermalization with partonic collectivity before expansion and freeze-out. It is expected that as
1218 the system size decreases and energy decreases, the condition for a complete thermalization are no
1219 longer met, and this change of condition may result in change in elliptic flow. Thus it becomes
1220 important to measure elliptic flow in smaller systems and at different energies. In this analysis,
1221 we present STAR's preliminary results of v_2 for K_S^0 , Φ , and Λ from 200 GeV CuCu collisions. We
1222 will compare the results with those from AuAu collisions at a similar number of participants or
1223 similar eccentricity. Hydrodynamic model results will also be used to compare with the data.

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Systematic Study of Low- p_T Direct Photons via π^0 Tagging Method at PHENIX

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1227 Direct photons are an excellent penetrating probe to study the formation of the Quark-Gluon
1228 Plasma (QGP) since they are not influenced by the strong interaction and hadronization
1229 processes. While at high p_T they constitute a crucial reference for the medium modification of
1230 high- p_T hadron production, at low and intermediate p_T direct photons, which are expected
1231 dominated by thermal photons, can provide important information about temperature of QGP at
1232 the early stage of heavy-ion collisions. But the huge background from meson decay makes the
1233 measurement of thermal photons very challenging.

1234 Among all the different efforts to measure low- p_T direct photons production at PHENIX, π^0
1235 tagging method is a promising approach to minimize the systematic errors. The analysis utilizes
1236 two different Electromagnetic Calorimeter (EMCal) detectors installed at PHENIX that provides
1237 an important cross checks and increases the confidence level of physics results. Methods, results
1238 and comparison with the latest measurements of PHENIX will be presented.

Three Body Force for investigation of Hadron in Hypercentral Harmonic and Anharmonic Potentials

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1239 We present a theoretical approach to the internal motion of a system based on 3-body forces
1240 among particles in a special case. Using 3-body potentials. The 3body-forces are more easily
1241 introduced and treated within the hyperspherical harmonic formalism. The internal particle
1242 motion is usually described by means of the Jacobin relative coordinates ρ , λ and R . The
1243 problems related to three body nonhypercentral potentials in three-dimension are investigated.

1244 While the difficulties that arise in the study of nonhypercentral potentials are explicitly shown,
1245 we discuss some results obtained using nonhypercentral harmonic and anharmonic and some
1246 inverse power terms; however the potential can be easily generalized in order to allow a
1247 systematic analysis. Which admit an exact solution of the wave equation. The method is also
1248 applied to some other types of potentials.

1249

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Di-electron spectroscopy in CBM

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1255 The Compressed Baryonic Matter (CBM) experiment at the future Facility for Antiproton and
1256 Ion Research (FAIR) at GSI in Darmstadt will be a dedicated heavy-ion experiment exploring the
1257 QCD phase diagram in the region of moderate temperatures and high baryon densities with A+A
1258 collisions at 10-45 AGeV beam energy. This region of the phase diagram attracted strong
1259 attention in recent years because several interesting features of QCD can be addressed
1260 experimentally. An experimental observation of a first order phase transition or the critical end
1261 point would be a milestone in our understanding of the quark-hadron phase diagram.

1262 One key aspect of the CBM physics program is to look for the signal of chiral symmetry
1263 restoration by studying the in-medium properties of low mass vector mesons created inside the
1264 fireball. The leptonic decay channels of these mesons are of special interest as the leptons leave
1265 the hot and dense fireball without further interaction and may reveal information on the
1266 characteristics of the matter created in the collisions. Along with this, the measurement of
1267 charmonium, also in its leptonic decay channel, will provide important insight into the subject of
1268 charm production at threshold, charmonium formation in hot and dense matter and
1269 deconfinement phase transition.

1270 The vector mesons can be reconstructed from di-electron or di-muon channel. Here we will
1271 present the feasibility studies investigating the di-electron channel, within the proposed CBM
1272 setup. The electrons are identified using information from the Ring Imaging Cherenkov (RICH)
1273 detector and the Transition Radiation Detectors (TRD). A further clean up of the identified
1274 electron sample is also done using the Time Of Flight (TOF) signal. Momentum information is
1275 provided from a set of Silicon Tracking Stations (STS) located inside a dipole magnetic field
1276 immediately behind the target. For the low mass vector mesons the special challenge is to reduce
1277 the combinatorial background coming from random combinations of electrons and positrons. In
1278 CBM this becomes further difficult because of the fact that electron identification is provided only
1279 after the magnetic field required for tracking. Several topological cuts have been developed to
1280 reduce the background coming from π^0 Dalitz decay and conversion of γ in various materials in
1281 the detector systems. Results from these feasibility studies will be presented and compared with
1282 those from existing experiments. Similar effort has been put forward to reconstruct J/ψ through
1283 the di-electron channel and results of this analysis will also be discussed.