

# TOWARDS A CONTROLLED STUDY OF THE QCD CRITICAL POINT

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The phase diagram of QCD, as a function of temperature  $T$  and baryon chemical potential  $\mu_B$ , may contain a critical point  $(\mu_E, T_E)$  whose non-perturbative nature makes it a natural object of lattice studies. However, the sign problem prevents the application of standard Monte Carlo techniques at non-zero baryon density. We have been pursuing an approach free of the sign problem, where the chemical potential is taken as imaginary and the results are Taylor-expanded in  $\mu/T$  about  $\mu = 0$ , then analytically continued to real  $\mu$  [1,2].

Within this approach we have determined the curvature of the pseudo-critical temperature  $dT_c/d\mu^2|_{\mu=0}$ , which we can compare to the experimentally determined freeze-out curve [3].

We have also determined the sensitivity of the critical chemical potential  $\mu_E$  to the quark mass,  $d(\mu_E)^2/dm_q|_{\mu_E=0}$ . Our study indicates that the critical point moves to *smaller* chemical potential as the quark mass *increases* [3,4]. This finding, contrary to common wisdom, implies that the deconfinement crossover, which takes place in QCD at  $\mu = 0$  when the temperature is raised, will remain a crossover in the  $\mu$ -region where our Taylor expansion can be trusted. If this result, obtained on a coarse lattice, is confirmed by simulations on finer lattices now in progress, then we predict that no *chiral* critical point will be found for  $\mu_B \lesssim 500$  MeV, unless the phase diagram contains additional transitions.

## References

- [1] P. de Forcrand and O. Philipsen, Nucl. Phys. B **642** (2002) 290 [arXiv:hep-lat/0205016].
- [2] P. de Forcrand and O. Philipsen, Nucl. Phys. B **673** (2003) 170 [arXiv:hep-lat/0307020].
- [3] P. de Forcrand and O. Philipsen, JHEP **0701** (2007) 077 [arXiv:hep-lat/0607017].
- [4] P. de Forcrand, S. Kim and O. Philipsen, PoS (LATTICE2007) 178, arXiv:0711.0262 [hep-lat].